

Comment on “Microscopic Theory of Network Glasses”

In a recent Letter, Hall and Wolynes [1] (HW) ask whether a microscopic theory of network glasses can be developed starting from a model of dense spherical fluids. To do so, they constrain the number of nearest neighbors and count their central force interactions separately. They obtain the dynamical transition temperature T_A (below which the system is nonergodic and the motion is landscape determined), and the entropy crisis (Kauzman) temperature T_K as functions of n_b , the average number of nearest neighbor bonds/atom. A Lindemann melting criterion on the amplitude of nearest neighbor vibrations defines the glass transition temperature T_G . The model shows that T_A/T_G and T_K/T_G monotonically increase and decrease, respectively, as functions of n_b . For $n_b = 2.4$, $T_A/T_G = 100$, an unreasonable result. Mode coupling theory defines a critical nonergodicity temperature T_c beyond which a radical change in the long time limit of the density-density correlation function occurs. T_c has been plausibly estimated for vitreous silica in molecular dynamics simulations [2] as $T_c/T_G = 2$. Even this temperature is presently outside the reach of experimental investigations [2].

How do these results compare with experiments? There are standard procedures for extrapolating specific heats to obtain T_K , but identifying the onset of nonergodicity at T_A (or T_c) is much more difficult. Chalcogenide glasses are ideal test systems because one can synthesize them over a wide range of n_b by chemical alloying group IV additives in Se base glass. Fortunately, glass transitions of these systems have been recently reinvestigated [3] using T -modulated scanning calorimetry (MDSC), a new method which permits separating the usual DSC heat flow endotherm \dot{H}_T into a reversing part \dot{H}_{rev} which is ergodic (and which follows the modulated T profile) from the nonreversing part \dot{H}_{nr} which is nonergodic (arising from underlying temperature dependent activated processes) as illustrated in Fig. 1(a). MDSC permits one to establish this temperature T_A at which dynamics become landscape dominated (i.e., in MDSC language dominated by a T -dependent \dot{H}_{nr}) in contrast to the linear response regime (i.e., at high temperatures when the heat flow is \dot{H}_{rev} dominated by a constant activation energy).

In binary $\text{Ge}_x\text{Se}_{1-x}$ glasses, observed variations in T_A/T_G and T_K/T_G as a function of $n_b = 2 + 2x$ are compared to HW predictions in Fig. 1(b). The $T_K/T_G(n_b)$ results were obtained from a Vogel-Fulcher analysis of viscosity measurements [4]. One finds the observed and predicted variations in $T_K/T_G(n_b)$ ratio to be in reasonable accord with each other showing a general reduction starting from a value of about 0.9 at $n_b = 2$ to a value of 0.6 near $n_b = 2.7$. Note, however, that the broad global minimum in the observed T_K/T_G ratio near $n_b = 2.4$ is not reproduced by the HW approach. More serious is the fact that ob-

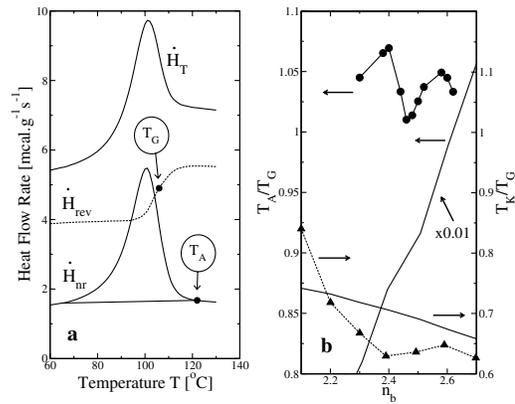


FIG. 1. (a) MDSC scan of $\text{As}_{20}\text{Se}_{80}$ bulk glass; T_G is defined as the inflexion point of \dot{H}_{rev} while T_A is the end point of the \dot{H}_{nr} endotherm. (b) T_A/T_G (●) and T_K/T_G (▲) for Ge-Se glasses. Lines without symbols are HW predictions [1].

served T_A/T_G values are (i) 2 orders of magnitude lower than the HW predictions and (ii) show a global minimum near $n_b = 2.4$ that is in sharp contrast to the almost linear increase [Fig. 1(b)] predicted by HW. The global minima in T_K/T_G and T_A/T_G ratios are features related to self-organization of glasses that are missing in the HW approach. Clearly, features such as inclusion of local structures [3,5], structural self-organization [5,6], and noncentral forces (bond bending) are missing in the theory. Bond-bending forces constitute $(4n_b - 6)/(5n_b - 6)$ of the global number of network constraints (e.g., 0.6 at the ideal $n_b = 2.4$) indicating that the noncentral (angular) forces have to be included in a successful theory of network glasses.

M. Micoulaut¹ and P. Boolchand²

¹LPTL, Université Paris VI
4 Place Jussieu 75252, Paris CEDEX 05, France
²Department of ECECS
University of Cincinnati
Cincinnati, Ohio 45221-0030, USA

Received 1 April 2003; published 9 October 2003

DOI: 10.1103/PhysRevLett.91.159601

PACS numbers: 61.43.Fs, 64.70.Pf, 65.60.+a

- [1] R. W. Hall and P. G. Wolynes, Phys. Rev. Lett. **90**, 085505 (2003).
- [2] J. Horbach and W. Kob, Phys. Rev. E **64**, 041503 (2001).
- [3] X. Feng, W. Bresser, and P. Boolchand, Phys. Rev. Lett. **78**, 4422 (1997); Y. Wang, J. Wells, D. G. Georgiev, P. Boolchand, K. Jackson, and M. Micoulaut, Phys. Rev. Lett. **87**, 185503 (2001).
- [4] C. A. Angell, in *Rigidity Theory and Applications*, edited by M. F. Thorpe and P. M. Duxbury (Plenum, New York, 1999).
- [5] M. Micoulaut and J. C. Phillips, Phys. Rev. B **67**, 104204 (2003); J. C. Phillips, Phys. Rev. Lett. **88**, 216401 (2002).
- [6] P. Boolchand, D. G. Georgiev, and B. Goodman, J. Optoelectron. Adv. Mater. **3**, 703 (2001).